Limits on a muon flux from Kaluza-Klein dark matter annihilations in the Sun from the IceCube 22-string detector


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A search for muon neutrinos from Kaluza-Klein dark matter annihilations in the Sun has been performed with the 22-string configuration of the IceCube neutrino detector using data collected in 104.3 days of live time in 2007. No excess over the expected atmospheric background has been observed. Upper limits have been obtained on the annihilation rate of captured lightest Kaluza-Klein particle (LKP) WIMPs in the Sun and converted to limits on the LKP-proton cross sections for LKP masses in the range 250–3000 GeV. These results are the most stringent limits to date on LKP annihilation in the Sun.

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In a recent work [1], we presented the result of a search for neutralino dark matter accumulated in the center of the Sun with the 22-string configuration of the IceCube detector. In this paper we extend the search to an alternative dark matter candidate, Kaluza-Klein (KK) particles, arising from theories with extra spacetime dimensions. In the simplest framework of universal extra dimensions (UED) [2], there is a single extra dimension of size $R \sim O(\text{TeV}^{-1})$ compactified on an $S^1/Z_2$ orbifold. Within minimal UED theories, the first excitation of the hypercharge gauge boson, $B^{(1)}$, is generally the lightest Kaluza-Klein particle (LKP). It is often denoted as the KK “photon,” $\gamma^{(1)}$, since the effective first KK-level Weinberg angle of the mass matrix is very small, and therefore $B^{(1)}$ can also be described as a mass eigenstate [2]. KK-parity conservation, affiliated with extra-dimensional momentum conservation, leads to the stability of the LKP, which makes it a viable dark matter (DM) candidate. There are also other possible natural choices for LKP candidates within UED, like the KK “graviton,” the KK “neutrino” or the $Z^{(1)}$ boson that may constitute viable DM candidates. They are not considered here. Instead, we focus on the most promising KKDM prospect in terms of indirect detection expectations, the KK “photon.” Accelerator measurements constrain the lower bound for the mass of the LKP, $m_{\gamma^{(1)}}$, at 300 GeV [3]. The upper bound is limited to a few TeV in order to not exceed the observed DM relic density and over close the Universe. Within this mass range, the relic density of LKPs is in the favorable region for providing the cold dark matter of the Universe, as observed by [4]. LKPs are assumed to have been in thermal equilibrium in the early Universe and therefore their relic density depends on the LKP mass, annihilation cross section and the coannihilation rate. The annihilation and coannihilation cross sections are determined by the mass spectrum of the first KK level and standard model couplings [5]. We here consider UED models with five spacetime dimensions characterized by two parameters: the LKP mass, $m_{\gamma^{(1)}}$, and the

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values of hilation channels of interest are given in Table I for two GeV–TeV range. The branching ratios for the LKP annihilation channels of interest are given in Table I for two mass splitting \( \Delta_{q_{10}} = (m_{q_{10}} - m_{q_{10}^{\gamma}})/m_{q_{10}} \), where \( m_{q_{10}} \) is the mass of the first KK quark excitation, as discussed in [2, 5–7]. As a possible dark matter component of the halo, LKPs can become gravitationally trapped in massive celestial bodies like the Sun, accumulating in the object’s core. Since the LKP is a boson, pairwise annihilation is dominated by s-wave processes, creating standard model particles whose decay chains produce neutrinos in the GeV–TeV range. The branching ratios for the LKP annihilation channels of interest are given in Table I for two values of \( \Delta_{q_{10}} \) [8]. The neutrinos may escape the Sun and reach Earth. The search presented here aims at detecting LKP annihilations indirectly by observing an excess of such high energy neutrinos from the direction of the Sun. Despite the existence of various limits on neutralino induced neutrino fluxes from the Sun [1, 9–12], no corresponding limits for LKP annihilations have been previously reported.

For the results presented here, we use the same data set, 104.3 days live time taken with the 22-string configuration of IceCube in 2007, and the same analysis cuts as presented in [1]. This is justified since the signature of the expected signal at the detector is very similar for the LKP and neutralinos, considering the hardest \( \chi^0 \)-annihilation channel into \( W^+W^- \). The neutrino spectrum from annihilations of a LKP of a given mass in the center of the Sun is considerably harder than that of a neutralino of the same mass. However, oscillations and energy losses of the neutrinos on their way out of the Sun, like neutral current scattering, absorption, and \( \nu_x \) regeneration, smear out the energy spectra in a way that makes them comparable at Earth. Figure 1 shows an example of how the resulting muon spectra at the detector compare for a selected choice of neutralino and LKP masses at 250 and 3000 GeV. The analysis strategy used in [1] is therefore already optimized for the search of KK dark matter.

We simulated LKP annihilations in the Sun using WIMPSIM [13] for LKP masses \( m_{q_{10}} = 250, 500, 700, 900, 1100, 1500, 3000 \) GeV. We used \( \Delta_{q_{10}} = 0 \) with annihilation branching ratios from Table I. Since \( \Delta_{q_{10}} > 0 \) results in an increased neutrino flux due to the importance of the contributions from the \( \tau^+\tau^- \) and the direct neutrino channels, the choice of \( \Delta_{q_{10}} = 0 \) leads to a conservative limit. The background in the search for neutrinos from the Sun comes from air showers created in cosmic ray interactions in the atmosphere. The showers cause downward going atmospheric muon events, triggering the detector at \( \sim 500 \) Hz, and atmospheric muon neutrino events, triggering at \( \sim 4 \) mHz. When the Sun is below the horizon, the neutrino signal can be distinguished from the atmospheric muon background by selecting events with upward-going reconstructed tracks.

Atmospheric muon and neutrino background events were also generated [14, 15]. The propagation of muons and photons in the ice was simulated [16, 17] taking measured ice properties into account [18].

The events had to pass several selection criteria as described in [1] in order to reduce the content of atmospheric muon events. As a compromise between signal efficiency and background rejection, it was required that more than half of the events in the final data sample were neutrino induced. The observables used describe the quality of the track reconstructions and the geometry and time evolution of the hit pattern in the detector, and they were required to be well reproduced in simulations. The event selection consisted first in a series of unidimensional cuts on the selected event variables, and a final step that used two support vector machines (SVM). The SVMs were trained with simulated signal, and a set of experimental data, recorded in December 2007 and not used in this analysis since the Sun was above the horizon, was taken as background. A final sample was then defined from a cut on the combined two SVM output values, \( Q_1 \times Q_2 \) (see Fig. 1 in [1]). The analysis was performed in a blind manner such that the azimuth of the Sun is unknown until the selection criteria were finalized.

<table>
<thead>
<tr>
<th>Channel</th>
<th>Branching ratio</th>
<th>( \Delta_{q_{10}} = 0 )</th>
<th>( \Delta_{q_{10}} = 0.14 )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( e^+e^- ), ( \mu^+\mu^- ), ( \tau^+\tau^- )</td>
<td>0.20</td>
<td>0.23</td>
<td></td>
</tr>
<tr>
<td>( u\bar{u}, c\bar{c}, t\bar{t} )</td>
<td>0.11</td>
<td>0.077</td>
<td></td>
</tr>
<tr>
<td>( dd, ss, bb )</td>
<td>0.007</td>
<td>0.005</td>
<td></td>
</tr>
<tr>
<td>( \nu_e\bar{\nu}<em>e, \nu</em>\mu\bar{\nu}<em>\mu, \nu</em>\tau\bar{\nu}_\tau )</td>
<td>0.012</td>
<td>0.014</td>
<td></td>
</tr>
<tr>
<td>( \Phi, \Phi^+ )</td>
<td>0.023</td>
<td>0.027</td>
<td></td>
</tr>
</tbody>
</table>

FIG. 1 (color online). Comparison of simulated muon spectra from LKP, \( \gamma(1) \), and neutralino, \( \chi^0 \), annihilations observed in IceCube, for two WIMP masses, 250 and 3000 GeV, representing the boundaries of the investigated LKP model space.
The systematic uncertainties on the effective volume, $V_{\text{eff}}$, defined as the equivalent detector volume with 100% selection efficiency, are the same as the ones calculated in the WIMP analysis in [1], and are dominated by the uncertainties in photon propagation in the ice and the absolute digital optical module efficiency. They range from $\pm 19\%$ for the highest $m_{\chi}$ to $\pm 26\%$ for the lowest $m_{\chi}$ [19]. From the final event selection of the signal simulation we additionally derive the effective area for muon neutrinos from the direction of the Sun as a function of neutrino energy, see Fig. 2. Also shown in the figure is the median angular error, the median of the angle between the reconstructed muon and the neutrino direction, $\Theta$. The result includes systematic uncertainties and is an average over the austral winter, during which the Sun is below the horizon.

For the LKP signal models we then calculated the effective volume and, based on the distribution of the reconstructed angle to the Sun $\Psi$, we constructed confidence intervals at the 90% confidence level using the method outlined in [1]: to evaluate the signal content in the final event sample, hypothesis testing was done based on $\Psi$, the angle between the reconstructed track and the direction of the Sun. From simulations we find $f_s(\Psi)$, the probability distribution of $\Psi$ for the signal. By randomizing the azimuth angle in the final event sample of experimental data, $f_b(\Psi)$, the equivalent probability distribution is found for background. Defining $\xi = \frac{\mu_s}{n_{\text{obs}}}$, from the number of signal events $\mu_s$ and the observed number of events $n_{\text{obs}}$, we form the combined probability density $f_\xi(\Psi) = \xi \cdot f_s(\Psi) + (1 - \xi) \cdot f_b(\Psi)$. Based on $\xi_{\text{best}}$, the non-negative signal content that maximizes the likelihood, we form the logarithm of the likelihood ratio $R(\xi) = \log(\prod_{i=1}^{n_{\text{obs}}} f_i(\Psi_i) / f_i(\Psi_i | \xi))$ [20]. Comparing this with a $R_{\text{test}}(\xi)$ distribution of a large number of pseudoexperiments with $n_{\text{obs}}$ events taken from $f_\xi(\Psi)$, we construct the confidence interval on $\xi$ at significance $\alpha$ as $R(\xi_{\lim}) = R_{\text{test}}^\alpha(\xi_{\lim})$, where $P(R_{\text{test}} > R_{\text{test}}(\xi_{\lim})) = 1 - \alpha$.

No excess of events from the Sun above the background expectation was found in the search ($\xi_{\text{best}} = 0$). The observed number of events as a function of the angle to the Sun, $\Psi$, is compared to the atmospheric background expectation in Fig. 3. From the upper limits on the number of signal events $\mu_s$, we calculate the limit on the neutrino to muon conversion rate $\Gamma_{\nu\mu} = \frac{\mu_s}{n_{\text{obs}}} t$, for the live time $t$.

Using the signal simulation [13], we can convert this rate to a limit on the LKP annihilation rate in the Sun, $\Gamma_\chi$, see Table II. Results from different experiments are commonly compared by calculating the limit on the muon flux above 1 GeV, $\Phi_\mu$, which is also given in Table II together with the sensitivity, $\bar{\Phi}_\mu$, the median limit obtained from simulations with no signal.

The flux limit is shown in Fig. 4 together with the theoretically allowed flux region, derived from Refs. [6,8] with the use of Dark SUSY [21]. We have here approximated the branching ratios for the regions of $\Delta_{g(1)} = 0.01$ and $\Delta_{g(1)} = 0.1$ with those of $\Delta_{g(1)} = 0$ and $\Delta_{g(1)} = 0.14$, respectively, as given in Table I. The limits on the annihilation rate can be converted into limits on the spin-dependent, $\sigma^{\text{SD}}$, and spin-independent, $\sigma^{\text{SI}}$, LKP-proton cross sections, allowing a comparison with the results of direct search experiments. Since capture in the
The results are shown in Table II. We assumed a local WIMP density of 0.1 cm$^{-3}$ and used the method described in Ref. [22] to perform the conversion.

Figure 5 shows the limits on the spin-dependent LKP-proton cross sections in the Sun including systematic errors (squares), the nonexcluded LKP mass regions (dashed), and the upper bound on $m_{\chi^{(0)}}$ derived from the overclosure limit for each individual LKP model [5], varies with different values of $\Delta m^{(i)}$ and increases remarkably for models with $\Delta m^{(i)} < 0.1$. The upper bound on $m_{\chi^{(0)}}$ is obtained with the 22-string configuration of IceCube compared with other bounds [24–26], and the KK model space. The theoretical model space (green area) is plotted for different predictions for the mass splitting $\Delta m^{(i)}$. The blue regions indicate the overlap regions with two different $\Omega_{\text{CDM}}$ intervals,

whereas the narrow dark blue region corresponds to the preferred WMAP 1σ region for cold dark matter. The upper bound on $m_{\chi^{(0)}}$, derived from the overclosure limit for each individual LKP model [5], varies with different values of $\Delta m^{(i)}$ and increases remarkably for models with $\Delta m^{(i)} < 0.1$. This is due to additional coannihilation effects, arising for degenerate LKP models [6].

In conclusion, we have presented the first limits on LKP annihilations in the Sun. We also derived the most stringent limits on the spin-dependent LKP-proton cross sections in the nonexcluded LKP mass regions (300 GeV < $m_{\chi^{(0)}}$ < 3 TeV), improving existing limits by more than 2 orders of magnitude.
magnitude and excluding some viable LKP models. The full IceCube detector with the DeepCore extension [27] is expected to test most LKP models within the allowed region for $0.05 < \Omega_{\text{CDM}} h^2 < 0.20$, shown in Fig. 5.

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